# THE QUANTUM MACMAHON MASTER THEOREM 

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#### Abstract

We state and prove a quantum-generalization of MacMahon's celebrated Master Theorem, and relate it to a quantum-generalization of the boson-fermion correspondence of Physics.


## 1. Introduction

1.1. MacMahon's Master Theorem. In this paper we state and prove a quantum-generalization of MacMahon's celebrated Master Theorem, conjectured by the first two authors. Our result was motivated by quantum topology. In addition to its potential importance in knot theory and quantum topology (explained in brief in the last section), this paper answers George Andrews's long-standing open problem [A] of finding a natural q-analog of MacMahon's Master Theorem.

Let us recall the original form of MacMahon's Master Theorem and some of its modern interpretations.
Consider a square matrix $A=\left(a_{i j}\right)$ of size $r$ with entries in some commutative ring. For $1 \leq i \leq r$, let $X_{i}:=\sum_{j=1}^{r} a_{i j} x_{j}$, (where $x_{i}$ 's are commuting variables) and for any vector ( $m_{1}, \ldots, m_{r}$ ) of non-negative integers let $G\left(m_{1}, \ldots, m_{r}\right)$ be the coefficient of $x_{1}^{m_{1}} x_{2}^{m_{r}} \ldots x_{r}^{m_{r}}$ in $\prod_{i=1}^{r} X_{i}^{m_{i}}$. MacMahon's Master Theorem is the following identity (see $[\mathrm{MM}]$ ):

$$
\begin{equation*}
\sum_{m_{1}, m_{2}, \ldots, m_{r}=0}^{\infty} G\left(m_{1}, \ldots, m_{r}\right)=1 / \operatorname{det}(I-A) \tag{1}
\end{equation*}
$$

There are several equivalent reformulations of MacMahon's Master Theorem; see for example [FZ] and references therein. Let us mention one, of importance to physics.

Given a matrix $A=\left(a_{i j}\right)$ of size $r$ with commuting entries which lie in a ring $\mathcal{R}$, and a nonnegative integer $n$, we can consider its symmetric and exterior powers $S^{n}(A)$ and $\Lambda^{n}(A)$, and their traces $\operatorname{tr} S^{n}(A)$ and $\operatorname{tr} \Lambda^{n}(A)$ respectively. Since

$$
\begin{aligned}
\operatorname{tr} S^{n}(A) & =\sum_{m_{1}+\ldots m_{r}=n} G\left(m_{1}, \ldots, m_{r}\right) \\
\operatorname{det}(I-t A) & =\sum_{n=0}^{\infty}(-1)^{n} \operatorname{tr} \Lambda^{n}(A) t^{n}
\end{aligned}
$$

the following identity

$$
\begin{equation*}
\frac{1}{\sum_{n=0}^{\infty}(-1)^{n} \operatorname{tr} \Lambda^{n}(A) t^{n}}=\sum_{n=0}^{\infty} \operatorname{tr} S^{n}(A) t^{n} \tag{2}
\end{equation*}
$$

in $\mathcal{R}[[t]]$ is equivalent to (1). In Physics (2) is called the boson-fermion correspondence, where bosons (resp. fermions) are commuting (resp. skew-commuting) particles corresponding to symmetric (resp. exterior) powers.

[^0]1.2. Quantum algebra, right-quantum matrices and quantum determinants. In $r$-dimensional quantum algebra we have $r$ indeterminate variables $x_{i}(1 \leq i \leq r)$, satisfying the commutation relations $x_{j} x_{i}=q x_{i} x_{j}$ for all $1 \leq i<j \leq r$. We also consider matrices $A=\left(a_{i j}\right)$ of $r^{2}$ indeterminates $a_{i j}, 1 \leq i, j \leq r$, that commute with the $x_{i}$ 's and such that for any 2 by 2 minor of ( $a_{i j}$ ), consisting of rows $i$ and $i^{\prime}$, and columns $j$ and $j^{\prime}$ (where $1 \leq i<i^{\prime} \leq r$, and $1 \leq j<j^{\prime} \leq r$ ), writing $a:=a_{i j}, b:=a_{i j^{\prime}}, c:=a_{i^{\prime} j}, d:=a_{i^{\prime} j^{\prime}}$, we have the commutation relations:
\[

$$
\begin{align*}
c a & =q a c, \quad(q \text {-commutation of the entries in a column) }  \tag{3}\\
d b & =q b d, \quad(q \text {-commutation of the entries in a column) }  \tag{4}\\
a d & =d a+q^{-1} c b-q b c \quad \text { (cross commutation relation). } \tag{5}
\end{align*}
$$
\]

We will call such matrices $A$ right-quantum matrices.
The quantum determinant, (first introduced in [FRT]) of any (not-necessarily right-quantum) $r$ by $r$ matrix $B=\left(b_{i j}\right)$ may be defined by

$$
\operatorname{det}_{q}(B):=\sum_{\pi \in S_{r}}(-q)^{-\operatorname{inv}(\pi)} b_{\pi_{1} 1} b_{\pi_{2} 2} \cdots b_{\pi_{r} r}
$$

where the sum ranges over the set of permutations, $S_{r}$, of $\{1, \ldots, r\}$, and for any of its members, $\pi, \operatorname{inv}(\pi)$ denotes the number of pairs $1 \leq i<j \leq r$ for which $\pi_{i}>\pi_{j}$.
1.3. A $q$-version of MacMahon's Master Theorem. We are now ready to state our quantum version of Macmahon's Master Theorem.

Theorem 1. (Quantum MacMahon Master Theorem) Fix a right-quantum matrix $A$ of size $r$. For $1 \leq i \leq r$, let $X_{i}:=\sum_{j=1}^{r} a_{i j} x_{j}$, and for any vector $\left(m_{1}, \ldots, m_{r}\right)$ of non-negative integers let $G\left(m_{1}, \ldots, m_{r}\right)$ be the coefficient of $x_{1}^{m_{1}} x_{2}^{m_{r}} \ldots x_{r}^{m_{r}}$ in $\prod_{i=1}^{r} X_{i}^{m_{i}}$. Let

$$
\operatorname{Ferm}(A)=\sum_{J \subset\{1, \ldots, r\}}(-1)^{|J|} \operatorname{det}_{q}\left(A_{J}\right)
$$

where the summation is over the set of all subsets $J$ of $\{1, \ldots, r\}$, and $A_{J}$ is the $J$ by $J$ submatrix of $A$, and

$$
\operatorname{Bos}(A)=\sum_{m_{1}, \ldots, m_{r}=0}^{\infty} G\left(m_{1}, \ldots, m_{r}\right)
$$

Then

$$
\operatorname{Bos}(A)=1 / \operatorname{Ferm}(A)
$$

When we specialize to $q=1$, Theorem 1 recovers Equation (2), which explains why our result is a $q$-version of the MacMahon Master Theorem. For a motivation of Theorem 1, see Section 3.

The above result is not only interesting from the combinatorial point of view, but it is also a key ingredient in a finite noncommutative formula for the colored Jones function of a knot. This will be explained in a subsequent publication, [HL].
1.4. Computer code. The results of the paper have been verified by computer code, written by the third author. Maple programs QuantumMACMAHON and qMM are available at: http://www.math.rutgers.edu/ ~zeilberg/. The former proves rigorously Theorem 1 for any fixed $r$.
1.5. Acknowledgement. The authors wish to thank the anonymous referee who pointed out an error in an earlier version of the paper, and Martin Loebl for enlightening conversations.

## 2. Proof

2.1. Some lemmas on operators. The proof will make crucial use of a calculus of difference operators, developed by the third author in $[\mathrm{Z1}]$. This calculus of difference operators predates the more advanced calculus of holonomic functions, developed by the third author in [Z2].

Difference operators act on discrete functions $F$, that is functions whose domain is $\mathbb{N}^{r}$. For example, consider the shift-operators $M_{i}$ and the multiplication operator $Q_{i}$ which act on a discrete function $F\left(m_{1}, \ldots, m_{r}\right)$ by

$$
\begin{aligned}
\left(M_{i} F\right)\left(m_{1}, \ldots, m_{r}\right) & :=F\left(m_{1}, \ldots, m_{i-1}, m_{i}+1, m_{i+1}, \ldots, m_{r}\right) \\
\left(Q_{i} F\right)\left(m_{1}, \ldots, m_{r}\right) & :=q^{m_{i}} F\left(m_{1}, \ldots, m_{r}\right)
\end{aligned}
$$

It is easily seen that

$$
M_{i} Q_{i}=q Q_{i} M_{i}
$$

Abbreviating $Q_{i}$ by $q^{m_{i}}$, we obtain that:

$$
\begin{equation*}
M_{i} q^{m_{i}}=q^{m_{i}+1} M_{i} \quad M_{i} q^{m_{j}}=q^{m_{j}} M_{i} \quad \text { for } \quad i \neq j \tag{6}
\end{equation*}
$$

Another example is the operator $\hat{x}_{i}$ which left multiplies $F$ by $x_{i}$. Notice that $\hat{x_{j}} \hat{x}_{i}=q \hat{x}_{i} \hat{x}_{j}$ for $j>i$. In the proof below, we will denote $\hat{x}_{i}$ by $x_{i}$. In that case, the identity $x_{j} x_{i}=q x_{i} x_{j}$ for $j>i$ holds in the quantum algebra, as well as in the algebra of operators.

Before embarking on the proof, we need the following readily-verified lemmas.
Lemma 2.1. (commuting $X_{i}$ with $X_{j}$ ) For $1 \leq i<j \leq r, X_{j} X_{i}=q X_{i} X_{j}$.
Lemma 2.2. (commuting $x_{i}$ with $X_{j}$ ) For each of the $a_{i j}$, define the operator $Q_{i j}$ acting on expressions $P$ involving $a_{i j}$ by $Q_{i j} P\left(a_{i j}\right):=P\left(q a_{i j}\right)$. Then, for any $1 \leq i, j \leq r$, and integer $m_{i}$ and any expression $F$

$$
x_{i}^{-m_{i}} X_{j} F=\left[\left(Q_{j 1}^{-1} Q_{j 2}^{-1} \cdots Q_{j, i-1}^{-1} Q_{j, i+1} \cdots Q_{j r}\right)^{m_{i}} X_{j}\right] x_{i}^{-m_{i}} F
$$

Lemma 2.3. (Column expansion with respect to the last column): Given an $r$ by $r$ matrix ( $a_{i j}$ ) (not necessarily quantum) let $A_{i}$ be the minor of the entry $a_{i r}$, i.e. the $r-1$ by $r-1$ matrix obtained by deleting the $i^{\text {th }}$ row and $r^{\text {th }}$ column. Then

$$
\operatorname{det}_{q}(A)=\sum_{i=1}^{r}(-q)^{i-r}\left(\operatorname{det}_{q} A_{i}\right) a_{i r}
$$

Lemma 2.4. If $A$ is a matrix that satisfies Equation (5) and $A^{\prime}$ denotes a matrix obtained by interchanging the $i$ and $j$ columns columns of $A$, then $\operatorname{det}_{q}\left(A^{\prime}\right)=(-q)^{-\operatorname{inv}(i j)} \operatorname{det}_{q}(A)$.
Proof. Suppose first that we interchange two adjacent colums $i$ and $j:=i+1$. Consider the involution of $S_{r}$ that sends a permutation $\pi$ to $\pi^{\prime}=\pi(i j)$. Given $\pi \in S_{r}$, let $(A ; \pi)=(-1)^{-\operatorname{inv}(\pi)} a_{\pi_{1} 1} \ldots a_{\pi_{r} r}$ denote the contribution of $\pi$ in $\operatorname{det}_{q}(A)$. Then, $\operatorname{det}_{q}(A)=\sum_{\pi}(A ; \pi)$. Equation (5) implies that

$$
(A ; \pi)+\left(A ; \pi^{\prime}\right)=(-q)\left(\left(A^{\prime} ; \pi\right)+\left(A^{\prime} ; \pi^{\prime}\right)\right)
$$

Summing over all permutations proves the result when $j=i+1$.
Observe that when $j=i+1$, the matrix $A^{\prime}$ is no longer right-quantum since it does not satisfy (5). However, the proof used only the fact that (5) holds for the $i$ and $i+1$ columns of $A$.

Thus, the proof can be iterated $\operatorname{inv}(i j)$ times to commute the $i$ and $j>i$ columns of $A$. The result follows.

Lemma 2.5. (Equal columns imply that $\operatorname{det}_{q}$ vanishes): Let $A$ be a right-quantum matrix. In the notation of Lemma 2.3, for all $j \neq r$,

$$
\sum_{i=1}^{r}(-q)^{i-r}\left(\operatorname{det}_{q} A_{i}\right) a_{i j}=0
$$

Proof. If $j=r-1$, it is easy to see that $q$-commutation along the entries in every column of $A$ imply that the sum vanishes.

If $j<r-1$, use Lemma 2.4 to reduce it to the case of $j=r-1$.

Remark 2.6. One can give an alternative proof of Lemmas 2.4 and 2.5 from the trivial 2 by 2 case and, by induction using the $q$-Laplace expansion of a $q$-determinant that is completely analogous to the classical case.
2.2. Proof of Theorem 1. The proof is a quantum-adaptation of the "operator-elimination" proof of MacMahon's Master Theorem given in [Z1]. Fix a right-quantum matrix $A$.

Observe that $G\left(m_{1}, \ldots, m_{r}\right)$ is the coefficient of $x_{1}^{0} \ldots x_{r}^{0}$ in

$$
H\left(m_{1}, \ldots, m_{r} ; x_{1}, \ldots, x_{r}\right):=x_{r}^{-m_{r}} \cdots x_{2}^{-m_{2}} x_{1}^{-m_{1}} \prod_{i=1}^{r} X_{i}^{m_{i}}
$$

We will think of $H$ as a discrete function, that is as a function of $\left(m_{1}, \ldots, m_{r}\right) \in \mathbb{N}^{r} . H$ takes values in the ring of noncommutative Laurrent polynomials in the $x_{i} \mathrm{~s}$, with coefficients in the ring generated by the entries of $A$, modulo the ideal given by (3)-(5).

Let's see how the shift operators $M_{i}$ acts on $H$. By definition,

$$
\begin{gathered}
M_{i} H\left(m_{1}, \ldots, m_{r} ; x_{1}, \ldots, x_{r}\right)= \\
x_{r}^{-m_{r}} \cdots x_{i+1}^{-m_{i+1}} x_{i}^{-m_{i}-1} x_{i-1}^{-m_{i-1}} \cdots x_{1}^{-m_{1}} X_{1}^{m_{1}} \cdots X_{i-1}^{m_{i-1}} X_{i}^{m_{i}+1} X_{i+1}^{m_{i+1}} \cdots X_{r}^{m_{r}}
\end{gathered}
$$

By moving $x_{i}{ }^{-1}$ to the front and $X_{i}$ in front of $X_{1}{ }^{m_{1}}$, and using Lemma 2.1 and $x_{j} x_{i}=q x_{i} x_{j}$, we have

$$
M_{i} H\left(m_{1}, \ldots, m_{r} ; x_{1}, \ldots, x_{r}\right)=q^{m_{r}+m_{r-1}+\cdots+m_{i+1}-m_{1}-m_{2}-\cdots-m_{i-1}} x_{i}^{-1}\left[x_{r}^{-m_{r}} \cdots x_{1}^{-m_{1}} X_{i}\right] X_{1}^{m_{1}} \cdots X_{r}^{m_{r}}
$$

By moving $X_{i}$ next to $x_{i}^{-1}$ and using Lemma 2.2 this equals to:

$$
\begin{gathered}
q^{m_{r}+m_{r-1}+\cdots+m_{i+1}-m_{1}-m_{2}-\cdots-m_{i-1}} x_{i}^{-1} \\
{\left[\left(Q_{i 2} \cdots Q_{i r}\right)^{m_{1}}\left(Q_{i 1}^{-1} Q_{i 3} \cdots Q_{i r}\right)^{m_{2}}\left(Q_{i 1}^{-1} Q_{i 2}^{-1} Q_{i 4} \cdots Q_{i r}\right)^{m_{3}} \cdots\left(Q_{i 1}^{-1} Q_{i 2}^{-1} \cdots Q_{i, r-1}^{-1}\right)^{m_{r}} X_{i}\right]} \\
x_{r}^{-m_{r}} \cdots x_{1}^{-m_{1}} X_{1}^{m_{1}} \cdots X_{r}^{m_{r}}
\end{gathered}
$$

which is equal to

$$
\begin{gathered}
q^{m_{r}+m_{r-1}+\cdots+m_{i+1}-m_{1}-m_{2}-\cdots-m_{i-1}} x_{i}^{-1} \\
\left(q^{-m_{2}-m_{3}-\cdots-m_{r}} a_{i 1} x_{1}+q^{m_{1}-m_{3}-\cdots-m_{r}} a_{i 2} x_{2}+\cdots+q^{m_{1}+m_{2}+\cdots+m_{r-1}} a_{i r} x_{r}\right) H\left(m_{1}, \ldots, m_{r} ; x_{1}, \ldots, x_{r}\right)
\end{gathered}
$$

Multiplying out and rearranging, we get that the discrete function $H\left(m_{1}, \ldots, m_{r} ; x_{1}, \ldots, x_{r}\right)$ is annihilated by the $r$ operators $(i=1,2, \ldots, r)$

$$
\mathcal{P}_{i}:=\sum_{j=1}^{i-1}-q^{-m_{j}-2 m_{j+1}-\cdots-2 m_{i-1}-m_{i}} a_{i j} x_{j}+\left(M_{i}-a_{i i}\right) x_{i}+\sum_{j=i+1}^{r}-q^{m_{i}+2 m_{i+1}+\cdots+2 m_{j-1}+m_{j}} a_{i j} x_{j}
$$

Now comes a nice surprise. Let us define $b_{i j}$ to be the coefficient of $x_{j}$ in $\mathcal{P}_{i}$. For example, for $r=3$ we have:

$$
B=\left(\begin{array}{ccc}
M_{1}-a_{11} & -q^{m_{1}+m_{2}} a_{12} & -q^{m_{1}+2 m_{2}+m_{3}} a_{13} \\
-q^{-m_{1}-m_{2}} a_{21} & M_{2}-a_{22} & -q^{m_{2}+m_{3}} a_{23} \\
-q^{-m_{1}-2 m_{2}-m_{3}} a_{31} & -q^{-m_{2}-m_{3}} a_{32} & M_{3}-a_{33}
\end{array}\right)
$$

Lemma 2.7. $B$ is a right-quantum matrix.
Proof. It is easy to see that the entries in each column of $B q$-commute. To prove Equation (5), consider the following cases for a 2 by 2 submatrix $C$ of $B$ : $C$ contains two, (resp. one, resp. no) diagonal entries of $B$, and prove it case by case, using the fact that the operators $M_{i}$ and $q^{m_{j}}$ commute with the $a_{i j}$, and satisfy the commutation relations (6).

Now we eliminate $x_{1}, x_{2}, \ldots, x_{r-1}$ by left-multiplying $\mathcal{P}_{i}$ by the minor of $b_{i r}$ in $B=\left(b_{i j}\right)$ times $(-q)^{i-r}$, for each $i=1,2, \ldots, r$, and adding them all up. Since $B$ is right-quantum (by Lemma 2.7), Lemma 2.5 implies that the coefficients of $x_{1}, \ldots, x_{r-1}$ all vanish, and $\operatorname{det}_{q}(B) x_{r} H=0$. After left multiplying by $x_{r}^{-1}$ which commutes with the entries in $B$, we obtain that

$$
\operatorname{det}_{q}(B) H\left(m_{1}, \ldots, m_{r} ; x_{1}, \ldots, x_{r}\right)=0
$$

Since the entries of $B$ do not contain $x_{i}$ 's, it follows that $\operatorname{det}_{q}(B)$ annihilates every coefficient of $H$, in particular its constant term. Taking the constant term yields

$$
\operatorname{det}_{q}(B) G\left(m_{1}, \ldots, m_{r}\right)=0
$$

Here comes the next surprise.
Lemma 2.8. (a) We have:

$$
\operatorname{det}_{q}(B)=\sum_{J \subset\{1, \ldots, r\}}(-1)^{|J|} \operatorname{det}_{q}\left(A_{J}\right) M_{\bar{J}}
$$

where $\bar{J}=\{1, \ldots, r\}-J$ and $M_{J}=\prod_{j \in J} M_{j}$.
(b) In particular,

$$
\left.\operatorname{det}_{q}(B)\right|_{M_{1}=\cdots=M_{r}=1}=\operatorname{Ferm}(A)
$$

Proof. Let us expand $\operatorname{det}_{q}(B)$ as a sum over permutations $\pi \in S_{r}$. We have:

$$
\begin{aligned}
\operatorname{det}_{q}(B) & =\sum_{\pi \in S_{r}}(-q)^{-\operatorname{inv}(\pi)} b_{\pi_{1} 1} b_{\pi_{2} 2} \cdots b_{\pi_{r} r} \\
& =\sum_{\pi \in S_{r}} \prod_{i=1}^{r}(-q)^{-\operatorname{inv}(\pi, i)} b_{\pi_{i} i}
\end{aligned}
$$

where $\operatorname{inv}(\pi, i)$ is the number of $j>i$ such that $\pi_{i}>\pi_{j}$. Now, $b_{i j}=\delta_{i j} M_{i}-q_{i j} a_{i j}$, where $q_{i j}$ is a monomial in the variables $q^{m_{k}}$, and $\prod_{i} q_{\pi_{i} i}=1$. Moreover, if $\pi_{i}=i$, then for each $j$ with $i<j \neq \pi_{j}$, the exponent of $q^{m_{i}}$ in $q_{i j}$ is 2 if $\pi_{j}<i$ and 0 if $\pi_{j}>i$.

Since $\prod_{i} q_{\pi_{i} i}=1$, we can move the monomials $q_{i j}$ in the left of $\prod_{i}(-q)^{-\operatorname{inv}(\pi, i)} b_{\pi_{i} i}$, and then cancel them. The monomials commute with all entries of the matrix $b_{i j}$, except with the diagonal ones. Commuting $q^{2 m_{i}}$ with $b_{i i}=\delta_{\pi_{i} i} M_{i i}-q_{\pi_{i} i} a_{\pi_{i} i}$ gives: $b_{i i} q^{2 m_{i}}=q^{2 m_{i}}\left(\delta_{\pi_{i} i} q^{2} M_{i i}-q_{\pi_{i} i} a_{\pi_{i} i}\right)$. In other words, it replaces $M_{i}$ by $q^{2} M_{i}$. Thus, we have:

$$
\begin{aligned}
\operatorname{det}_{q}(B) & =\sum_{\pi \in S_{r}} \prod_{i=1}^{r}(-q)^{-\operatorname{inv}(\pi, i)}\left(\delta_{\pi_{i} i} q^{2 \operatorname{inv}(\pi, i)} M_{i}-a_{\pi_{i} i}\right) \\
& =\sum_{\pi \in S_{r}} \sum_{J \subset\{1, \ldots, r\}} \prod_{i \in J}(-q)^{-\operatorname{inv}(\pi, i)} \delta_{\pi_{i} i} q^{2 \operatorname{inv}(\pi, i)} M_{i} \prod_{i \notin J}(-q)^{-\operatorname{inv}(\pi, i)}\left(-a_{\pi_{i} i}\right)
\end{aligned}
$$

Now, rearrange the summation. Observe that every permutation $\pi$ of $\{1, \ldots, r\}$ gives rise to a permutation $\pi^{\prime}$ on the set $\{1, \ldots, r\}-\operatorname{Fix}(\pi)$, where $\operatorname{Fix}(\pi)$ is the fixed point set of $\pi$. Moreover, $\operatorname{inv}\left(\pi^{\prime}, i\right)=\operatorname{inv}(\pi, i)-\mid\{j \in$ $J: j>i\} \mid$. Using this, part (a) follows. Part (b) follows from part (a) and the definition of Ferm $(A)$.

Hence

$$
\sum_{J \subset\{1, \ldots, r\}}(-1)^{|J|} \operatorname{det}_{q}\left(A_{J}\right) M_{\bar{J}} G\left(m_{1}, \ldots, m_{r}\right)=0 .
$$

Summing over $\mathbb{N}^{r}$, we get:

$$
\sum_{m_{1}, \ldots, m_{r}=0}^{\infty} \sum_{J \subset\{1, \ldots, r\}}(-1)^{|J|} \operatorname{det}_{q}\left(A_{J}\right) M_{\bar{J}} G\left(m_{1}, \ldots, m_{r}\right)=0
$$

For a subset $J=\left\{k_{1}, \ldots, k_{j}\right\}$ of $\{1, \ldots, r\}$, we denote by $G_{J}\left(m_{k_{1}}, \ldots, m_{k_{j}}\right)$ the evaluation $G\left(m_{1}, \ldots, m_{r}\right)$ at $m_{i}=0$ for all $i \notin J$, and we define

$$
S_{J}=\sum_{m_{k_{1}}, \ldots, m_{k_{j}}=0}^{\infty} G\left(m_{1}, \ldots, m_{r}\right) .
$$

Using telescoping cancellation, the inclusion-exclusion principle, and Lemma 2.8(b), the above equation becomes

$$
\sum_{J \subset\{1, \ldots, r\}}(-1)^{|J|} \operatorname{Ferm}\left(A_{J}\right) S_{J}=0 .
$$

Using induction (with respect to $r$ ), together with $S_{\emptyset}=1$, we obtain that $\operatorname{Ferm}(A) S_{\{1, \ldots, r\}}=1$. This concludes the proof of the theorem.

## 3. Some remarks on the boson-FERMION CORRESPONDENCE

Let us give some motivation for Theorem 1 from the point of view of quantum topology.
For a reference on quantum space and quantum algebra, see [Ka, Chapter IV] and [M].
Recall that a vector (column or row) of $r$ indeterminate entries $x_{1}, \ldots, x_{r}$ lies in $r$-dimensional quantum space $A^{r \mid 0}$ iff its entries satisfy

$$
x_{j} x_{i}=q x_{i} x_{j}
$$

for all $1 \leq i<j \leq r$.
Recall that a right (resp. left) endomorphism of $A^{r \mid 0}$ is a matrix $A=\left(a_{i j}\right)$ of size $r$ whose entries commute with the coordinates $x_{i}$ of a vector $x=\left(x_{1}, \ldots, x_{r}\right)^{T} \in A^{r \mid 0}$ and in addition, $A x$ (resp. $x^{T} A$ ) lie in $A^{r \mid 0}$. Recall also that an endomorphism of $A^{r \mid 0}$ is one that is right and left endomorphism.

It is easy to see (eg. in [Ka, Thm. IV.3.1]) that $A$ is a right-quantum (i.e., a right-endomorphism) iff for every 2 by 2 submatrix $\left(\begin{array}{ll}a & b \\ c & d\end{array}\right)$ of $A$ we have:

$$
c a=q a c, \quad d b=q b d, \quad a d=d a+q^{-1} c b-q b c .
$$

Moreover, $A$ is left-quantum iff for every 2 by 2 submatrix of $A$ (as above) we have:

$$
b a=q a b, \quad d c=q c d, \quad a d=d a+q^{-1} b c-q c b
$$

Finally, $A$ is quantum iff for every 2 by 2 submatrix of $A$ (as above) we have:

$$
\begin{equation*}
b a=q a b, \quad c a=q a c, \quad d b=q b d, \quad d c=q c d, \quad c b=b c, \quad a d=d a+q^{-1} c b-q b c . \tag{7}
\end{equation*}
$$

The set of quantum matrices $A$ are the points of the $r$-dimensional quantum algebra $M_{q}(r)$, which is defined to be the quotient of the free algebra in noncommuting variables $x_{i j}$ for $1 \leq i, j, \leq r$, modulo the left ideal generated by the commutation relations of Equation (7).

The algebra $M_{q}(r)$ has interesting and important structure. $M_{q}(r)$ is Noetherian, has no zero divisors, and in addition, a basis for the underlying vector space is given by the set of sorted monomials $\left\{\prod_{i, j} a_{i j}^{n_{i j}} \mid n_{i j} \geq 0\right\}$ where the product is taken lexicographically; see [Ka, Thm IV.4.1]. An important quotient of $M_{q}(r)$ is the quantum group $S L_{q}(r):=M_{q}(r) /\left(\operatorname{det}_{q}-1\right)$, which is a Hopf algebra [Ka, Sec.IV.6] whose representation theory gives rise to the quantum group invariants of knots, such as the celebrated Jones polynomial.

Observing that

$$
\begin{aligned}
\operatorname{tr} S^{n}(A) & =\sum_{m_{1}+\ldots m_{r}=n} G\left(m_{1}, \ldots, m_{r}\right) \\
\operatorname{tr} \Lambda^{n}(A)= & \sum_{J \subset\{1, \ldots, r\},|J|=n} \operatorname{det}_{q}\left(A_{J}\right)
\end{aligned}
$$

Theorem 1 implies that:
Theorem 2. If $A$ is in $M_{q}(r)$, then

$$
\frac{1}{\operatorname{Ferm}(A)}=\sum_{n=0}^{\infty} \operatorname{tr} S^{n}(A)
$$

Since the algebra $M_{q}(r)$ has a vector space basis given by sorted monomials, it should be possible to give an alternative proof of the quantum MacMahon Master Theorem using combinatorics on words, as was done in [FZ] for several proofs of the MacMahon Master theorem. We hope to return to this alternative point of view in the near future.

## References

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