

2.5 n 10

We have

$$(i) \quad \partial_{x_i} u_\lambda(x, t) = \lambda \frac{\partial}{\partial \lambda x_i} u(\lambda x, \lambda^2 t)$$

$$\text{and } \partial_{x_i}^2 u_\lambda(x, t) = \lambda^2 \frac{\partial^2}{\partial (\lambda x_i)^2} u(\lambda x, \lambda^2 t)$$

while

$$\partial_t u_\lambda(x, t) = \lambda^2 \frac{\partial}{\partial (\lambda^2 t)} u(\lambda x, \lambda^2 t)$$

Thus

$$\partial_t u_\lambda(x, t) - \Delta u_\lambda(x, t) = \lambda^2 \left(\frac{\partial}{\partial z} u(y, z) - \Delta_y u(y, z) \right) = 0$$

where $z = \lambda^2 t$ and $y = \lambda x$.

(ii) Clearly we have

$$0 = \partial_\lambda \left(\partial_t u_\lambda(x, t) - \Delta u_\lambda(x, t) \right) = \partial_t \left(\partial_\lambda u_\lambda(x, t) \right) - \Delta \left(\partial_\lambda u_\lambda(x, t) \right)$$

so that $\partial_\lambda u_\lambda(x, t)$ solves the heat equation.

Observe that

$$v(x, t) = \partial_\lambda u_\lambda(x, t) \Big|_{\lambda=1}$$

2.5 n 12

Define

$$u(x, t) = e^{-ct} v(x, t).$$

Subst. to Ting we get

$$-ce^{-ct} v(x, t) + e^{-ct} \frac{\partial}{\partial t} v(x, t) - e^{-ct} \Delta v(x, t) + ce^{-ct} v(x, t) = f$$

or

$$\begin{cases} \frac{\partial}{\partial t} v(x, t) - \Delta v(x, t) = e^{ct} f(x, t) & \text{on } \mathbb{R}^n \times (0, \infty) \\ v(x, 0) = g(x) \end{cases}$$

Thus

$$v(x, t) = \int_{\mathbb{R}^n} \Phi(x-y, t) g(y) dy + \int_0^t \int_{\mathbb{R}^n} \Phi(x-y, t-s) e^{cs} f(y, s) dy ds$$

so that

$$u(x, t) = \int_{\mathbb{R}^n} e^{-ct} \Phi(x-y, t) g(y) dy + \int_0^t \int_{\mathbb{R}^n} \Phi(x-y, t-s) e^{-c(t-s)} f(y, s) dy ds$$

2.5 n 13

Consider

$$v(x,t) = u(x,t) - g(t)$$

IT solves

$$v_t - v_{xx} = -g'(t) \quad \text{on } \mathbb{R}^+ \times (0, \infty)$$

$$v = 0 \quad \text{on } \mathbb{R}^+ \times \{\tau = 0\}$$

$$v = 0 \quad \text{on } \{x = 0\} \times [0, \infty)$$

Let now

$$w(x,t) = \begin{cases} v(x,t) & x > 0 \\ -v(-x,t) & x < 0 \end{cases}$$

Reasoning as in ex 9 we find that

w satisfies

$$\begin{cases} w_t - w_{xx} = h(x,t) & \text{on } \mathbb{R}^+ \times (0, \infty) \\ w = 0 & \text{on } \mathbb{R} \times \{\tau = 0\} \end{cases}$$

so that

$$w(x,t) = \int_{\mathbb{R}} \phi(x-y, t-\tau) g'(\tau) dy + \int_0^t \int_{\mathbb{R}} \phi(x-y, t-\tau) h(y, \tau) dy d\tau$$

where

$$h(x,t) = \begin{cases} -g'(t) & x > 0 \\ +g'(t) & x < 0 \end{cases}$$

We thus have

$$\begin{aligned}
 w(x,t) &= \int_{\mathbb{R}^2} \int_0^t \phi(x-y, t-s) g'(s) dy ds = \\
 &= - \int_{\mathbb{R}^2} \int_0^t (\phi(x-y, t-s) - \phi(x+y, t-s)) g'(s) dy ds
 \end{aligned}$$

Integrating by part we get

$$\begin{aligned}
 w(x,t) &= - \int_{\mathbb{R}^2} \left[\frac{d}{dt} (\phi(x-y, t-s) - \phi(x+y, t-s)) \right] g(s) ds dy \\
 &\quad - g(t)
 \end{aligned}$$

where we used that

$$\lim_{\epsilon \rightarrow 0} \int_{\mathbb{R}^2} (\phi(x-y, \epsilon) - \phi(x+y, \epsilon)) dy = 1$$

Thus we have

$$w(x,t) = -g(t) - \int_{\mathbb{R}^2} \int_0^t \Delta_y (\phi(x-y, t-s) - \phi(x+y, t-s)) g(s) ds dy$$

~~Integrating by part~~ ^{Thus} we get

$$\begin{aligned}
 w(x,t) &= -g(t) - \int_0^t \left[\partial_y (\phi(x-y, t-s) - \phi(x+y, t-s)) \right] \Big|_{y=0}^{y=\infty} g(s) ds = \\
 &= -g(t) - 2 \int_0^t \partial_x \phi(x, t-s) g(s) ds
 \end{aligned}$$

So that we finally get, for $x > 0$

$$u(x, t) = w(x, t) + g(ct) =$$

$$- 2 \int \partial_x \phi(x, t-s) g(cs) ds$$

That is The desired formula.

2.5 n 14

(a) Repeating the proof of Theorem 3 page 52 we see that

$$\phi'(r) = \frac{1}{r^{n+1}} \iint_{E(r)} \left(-4n u_s \eta - \frac{2n}{s} \sum_{i=1}^n u_{y_i} y_i \right) dy ds$$

is valid for every $u(x,t)$ smooth.

We thus get

$$\phi'(r) \geq \frac{1}{r^{n+1}} \iint_{E(r)} \left(-4n \Delta u \eta - \frac{2n}{s} \sum_{i=1}^n u_{y_i} y_i \right) dy ds$$

Proceeding as on page 54 we get

$$\phi'(r) \geq 0$$

That, together with $\phi(0) = 4u(0,0)$

gives

$$\phi(r) \geq 4u(0,0)$$

That is exactly the Theorem

(b) Suppose There exists $(x_0, t_0) \in U$

such That

$$v(x_0, t_0) = \max_{\bar{U}_T} v$$

Find r such that $E(x_0, t_0; r) \subset U_T$.

We get

$$v(x_0, t_0) \leq \frac{1}{4r^n} \iint_{E(x_0, t_0; r)} v(y, s) \frac{|x_0 - y|^2}{|t - s|^2} dy ds$$

But since

$$\frac{1}{4r^n} \iint_{E(x_0, t_0; r)} \frac{|x_0 - y|^2}{|t - s|^2} dy ds = 1$$

we get

$$\frac{1}{4r^n} \iint_{E(x_0, t_0; r)} v(y, s) \frac{|x_0 - y|^2}{|t - s|^2} dy ds \leq \max_{\bar{U}_T} v = v(x_0, t_0)$$

So that

$$v(x_0, t_0) = \frac{1}{4r^n} \iint_{E(x_0, t_0; r)} v(y, s) \frac{|x_0 - y|^2}{|t - s|^2} dy ds$$

From here you can proceed exactly
as in the proof of Th. 4 pag. 54

(c) We get

$$\partial_t \phi(u) = \phi'(u) u_t$$

while

$$\Delta \phi(u) = \phi''(u) Du \cdot Du + \phi'(u) \Delta u$$

So that

$$v_t - \Delta v = -\phi''(u) |Du|^2 \leq 0$$

(d) Observe that if u solve the heat eq. so do u_t and u_{x_i} .

Since $\phi(x) = x^2$ is convex we have

that

$$u_t^2 \text{ and } u_{x_i}^2$$

are sub solution. ~~But~~ But

clearly the sum of sub solution is

a sub solution.