

TILINGS

APERIODIC MEDIA

and their

NONCOMMUTATIVE GEOMETRY

Jean BELLISSARD ^{1 2}

Georgia Institute of Technology

&

Institut Universitaire de France

Collaborations:

D. SPEHNER (Essen, Germany)

M. ZARROUATI (IRSAMC, Toulouse)

D. HERRMANN (T.U. Berlin)

J. KELLENDONK (U. Cardiff, United Kingdom)

A. LEGRAND (U. Paul Sabatier, Toulouse)

J.M. GAMBAUDO (U. Dijon, France)

R. BENEDETTI (U. Pisa, Italy)

¹Georgia Institute of Technology, School of Mathematics, Atlanta GA 30332-0160

²e-mail: jeanbel@math.gatech.edu

Main References

A. CONNES, *Sur la théorie non commutative de l'intégration*,
Lecture Notes in Math **725**, 19-143, Springer, Berlin (1979).

J. BELLISSARD, *The Gap Labelling Theorems for Schrödinger's Operators*,
in *From Number Theory to Physics*, pp. 538-630, Les Houches March 89, Springer,
J.M. Luck, P. Moussa & M. Waldschmidt Eds., (1993).

J.C. LAGARIAS, P.A.B. PLEASANT, *Repetitive Delone sets and perfect quasicrystals*,
in math.DS/9909033 (1999).

J. BELLISSARD, D. HERRMANN, M. ZARROUATI, *Hull of Aperiodic Solids and Gap Labelling Theorems*,
In *Directions in Mathematical Quasicrystals*, CRM Monograph Series,
Volume **13**, (2000), 207-259, M.B. Baake & R.V. Moody Eds., AMS Providence.

J. BELLISSARD, J. KELLENDONK, A. LEGRAND, *Gap Labelling for three dimensional aperiodic solids* ,
C. R. Acad. Sci. (Paris), **t.332**, Série I, p. 521-525, (2001).

J. BELLISSARD, R. BENEDETTI, J. M. GAMBAUDO, *Spaces of Tilings, Finite Telescopic Approximations,
and Gap-Labelling*, math.DS/0109062, (2001).

J. BELLISSARD, *Noncommutative Geometry of Aperiodic Solids*,
Proc. 2001 Summer School of Theor. Phys., Kluwer. *Geometry, Topology and Quantum Field Theory*,
Villa de Leyva, Colombia, 7-30 July 2001. <http://www.math.gatech.edu/~jeanbel/publications>

Content

1. The Hull as a Dynamical System
2. Building Hulls
3. Tilings & Point Sets
4. The Noncommutative Brillouin Zone
5. Gap labeling and K-theory.

Aperiodic Solids

1. *Perfect crystals* in d -dimensions:
translation and crystal symmetries.
Translation group $\mathcal{T} \simeq \mathbb{Z}^d$.
2. *Quasicrystals*: no translation symmetry, but icosahedral symmetry. Ex.:
 - (a) **Al_{62.5}Cu₂₅Fe_{12.5}**;
 - (b) **Al₇₀Pd₂₂Mn₈**;
 - (c) **Al₇₀Pd₂₂Re₈**;
3. *Amorphous media*: short range order
 - (a) Glasses;
 - (b) Silicon in amorphous phase;
4. *Disordered media*: random atomic positions
 - (a) Normal metals (with defects or impurities);
 - (b) Doped semiconductors (**Si**, **AsGa**, ...);

I - The Hull as a Dynamical System

J. BELLISSARD, D. HERMMANN, M. ZARROUATI, *Hull of Aperiodic Solids and Gap Labelling Theorems*
To appear in *Directions in Mathematical Quasicrystals*, M.B. Baake & R.V. Moody Eds, AMS, (2000).

I.1)- Point Sets

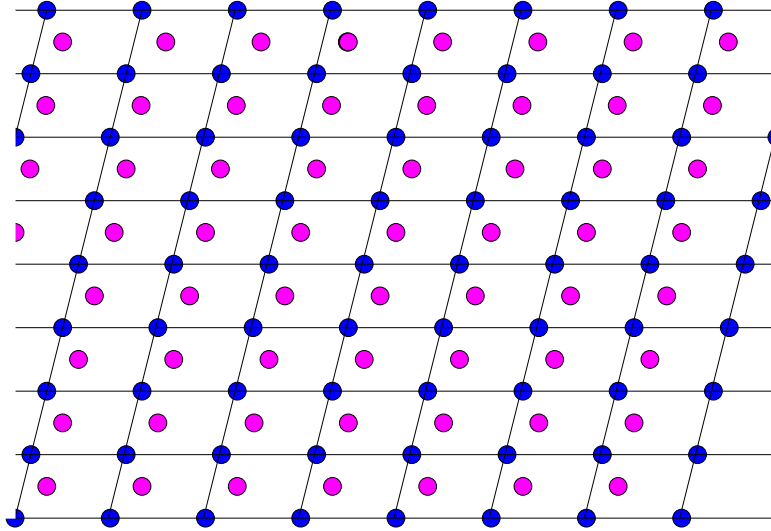
Equilibrium positions of atomic nuclei make up a point set $\mathcal{L} \subset \mathbb{R}^d$ *the set of lattice sites*. \mathcal{L} may be:

1. *Discrete*.
2. *Uniformly discrete*: $\exists r > 0$ s.t. each ball of radius r contains at most one point of \mathcal{L} .
3. *Relatively dense*: $\exists R > 0$ s.t. each ball of radius R contains at least one points of \mathcal{L} .
4. A *Delone* set: \mathcal{L} is uniformly discrete and relatively dense.
5. *Finite type Delone* set: $\mathcal{L} - \mathcal{L}$ is discrete.
6. *Meyer* set: \mathcal{L} and $\mathcal{L} - \mathcal{L}$ are Delone.

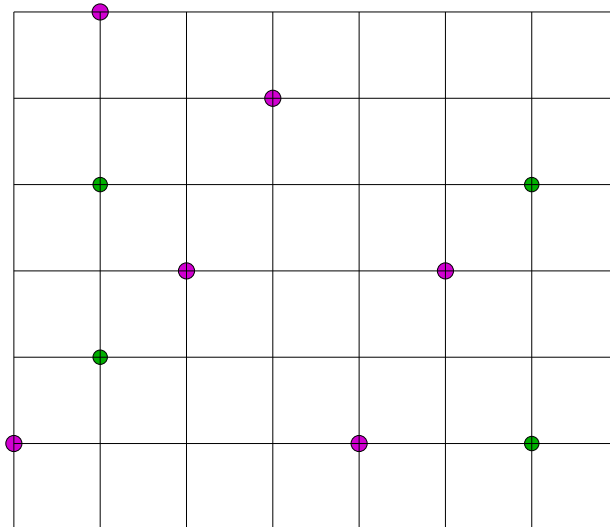
Examples:

1. A random Poissonian set in \mathbb{R}^d is almost surely discrete but not uniformly discrete nor relatively dense.
2. Due to Coulomb repulsion and Quantum Mechanics, **lattices of atoms are always uniformly discrete**.
3. Impurities in semiconductors are not relatively dense.
4. In amorphous media \mathcal{L} is Delone.
5. In a quasicrystal \mathcal{L} is Meyer.

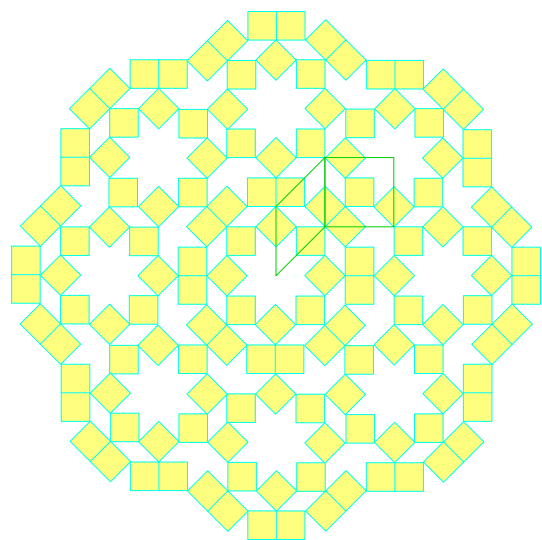
The ideal equilibrium positions of atomic nuclei sit on a discrete subset \mathcal{L} of \mathbb{R}^d ($d = 1, 2, 3$ in practice).



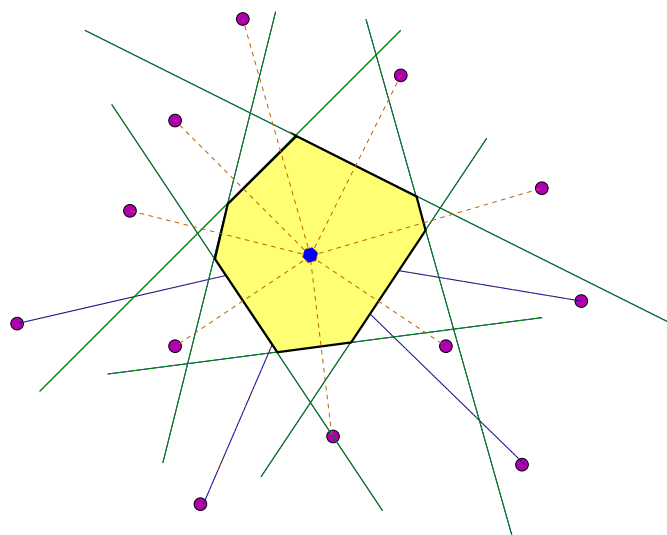
- *A periodic array of atomic nuclei* -



- *A random array of atomic nuclei* -



- *A quasiperiodic array of atomic nuclei* -



- *Construction of VORONOI's tiling* -

I.2)- Point Measures

$\mathfrak{M}(\mathbb{R}^d)$ is the set of Radon measures on \mathbb{R}^d namely the dual space to $\mathcal{C}_c(\mathbb{R}^d)$ (continuous functions with compact support), endowed with the weak* topology.

For \mathcal{L} a *uniformly discrete* point set in \mathbb{R}^d :

$$\nu := \nu^{\mathcal{L}} = \sum_{y \in \mathcal{L}} \delta(x - y) \in \mathfrak{M}(\mathbb{R}^d).$$

The *Hull* is the closure in $\mathfrak{M}(\mathbb{R}^d)$

$$\Omega = \overline{\{T^a \nu^{\mathcal{L}}; a \in \mathbb{R}^d\}},$$

where $T^a \nu$ is the translated of ν by a .

Results:

1. Ω is compact and \mathbb{R}^d acts by homeomorphisms.
2. If $\omega \in \Omega$, there is a uniformly discrete point set \mathcal{L}_ω in \mathbb{R}^d such that ω coincides with $\nu_\omega = \nu^{\mathcal{L}_\omega}$.
3. If \mathcal{L} is *Delone* (resp. *Meyer*) so are the \mathcal{L}_ω 's.

I.3)- Properties

(a) Minimality

\mathcal{L} is *repetitive* if for any finite patch p there is $R > 0$ such that each ball of radius R contains an ϵ -approximant of a translated of p .

Proposition 1 \mathbb{R}^d acts minimally on Ω if and only if \mathcal{L} is repetitive.

(b) Transversal

The closed subset $X = \{\omega \in \Omega ; \nu_\omega(\{0\}) = 1\}$ is called the *canonical transversal*. Let G be the subgroupoid of $\Omega \rtimes \mathbb{R}^d$ induced by X .

A Delone set \mathcal{L} has *finite type* if $\mathcal{L} - \mathcal{L}$ is closed and discrete.

(c) Cantorian Transversal

Proposition 2 If \mathcal{L} has finite type, then the transversal is completely discontinuous (Cantor).

II - Building Hulls

J. BELLISSARD, R. BENEDETTI, J.-M. GAMBAUDO, *Spaces of Tilings, Finite Telescopic Approximations and Gap-Labeling*, preprint August (2001).

II.1)- Examples

1. *Crystals* : $\Omega = \mathbb{R}^d / \mathcal{T} \simeq \mathbb{T}^d$ with the quotient action of \mathbb{R}^d on itself. (Here \mathcal{T} is the translation group leaving the lattice invariant. \mathcal{T} is isomorphic to \mathbb{Z}^D .)

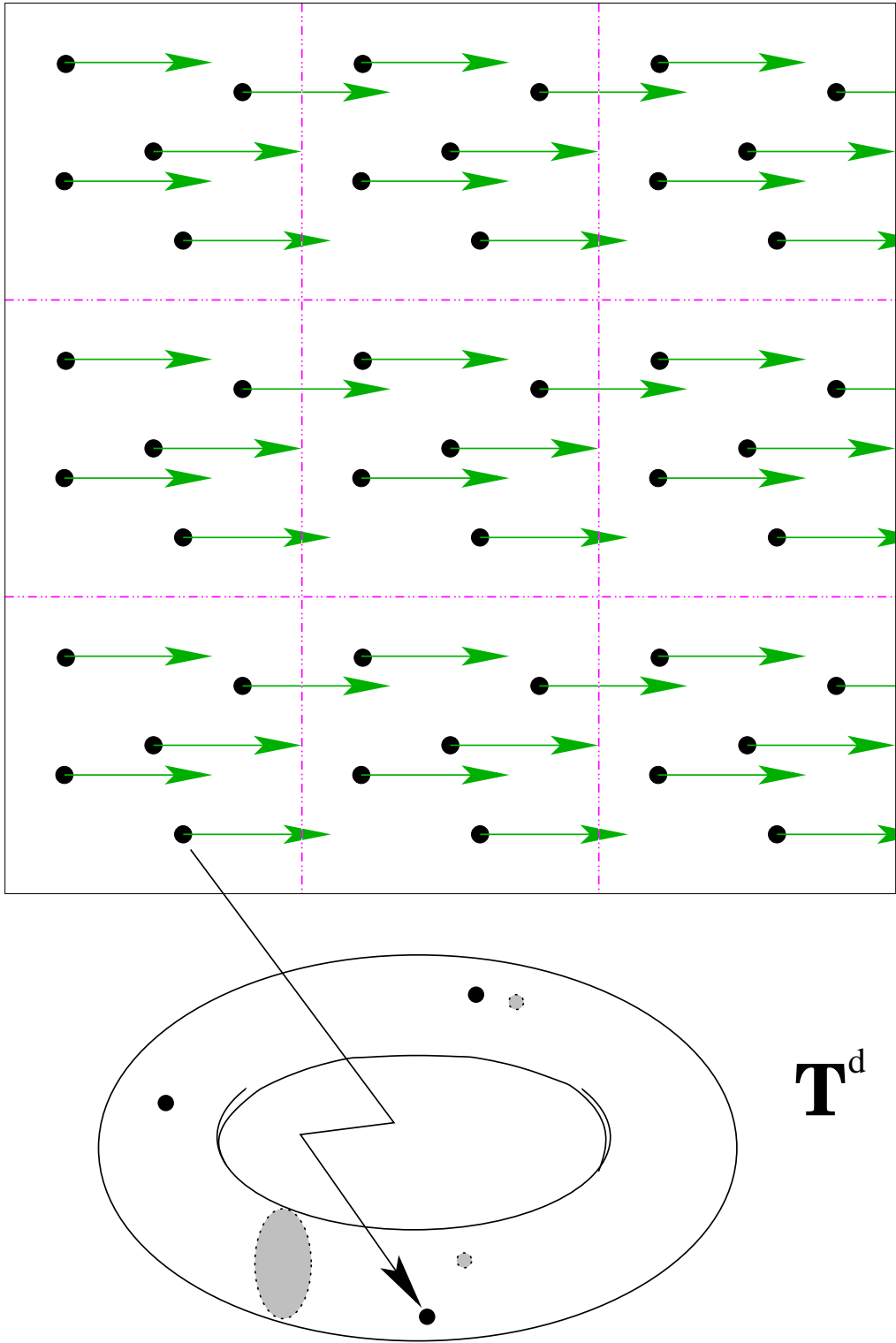
The transversal is a finite set (number of point per unit cell).

2. *Quasicrystals* : $\Omega \simeq \mathbb{T}^n$, $n > d$ with an irrational action of \mathbb{R}^d and a completely discontinuous topology in the transverse direction to the \mathbb{R}^d -orbits. The transversal is a Cantor set.

3. *Impurities in Si* : let \mathcal{L} be the lattices sites for *Si* atoms (it is a Bravais lattice). Let \mathfrak{A} be a finite set (alphabet) indexing the types of impurities.

The transversal is $X = \mathfrak{A}^{\mathbb{Z}^d}$ with \mathbb{Z}^d -action given by shifts.

The Hull Ω is the mapping torus of X .



- The Hull of a Periodic Lattice -

II.2)- Quasicrystals

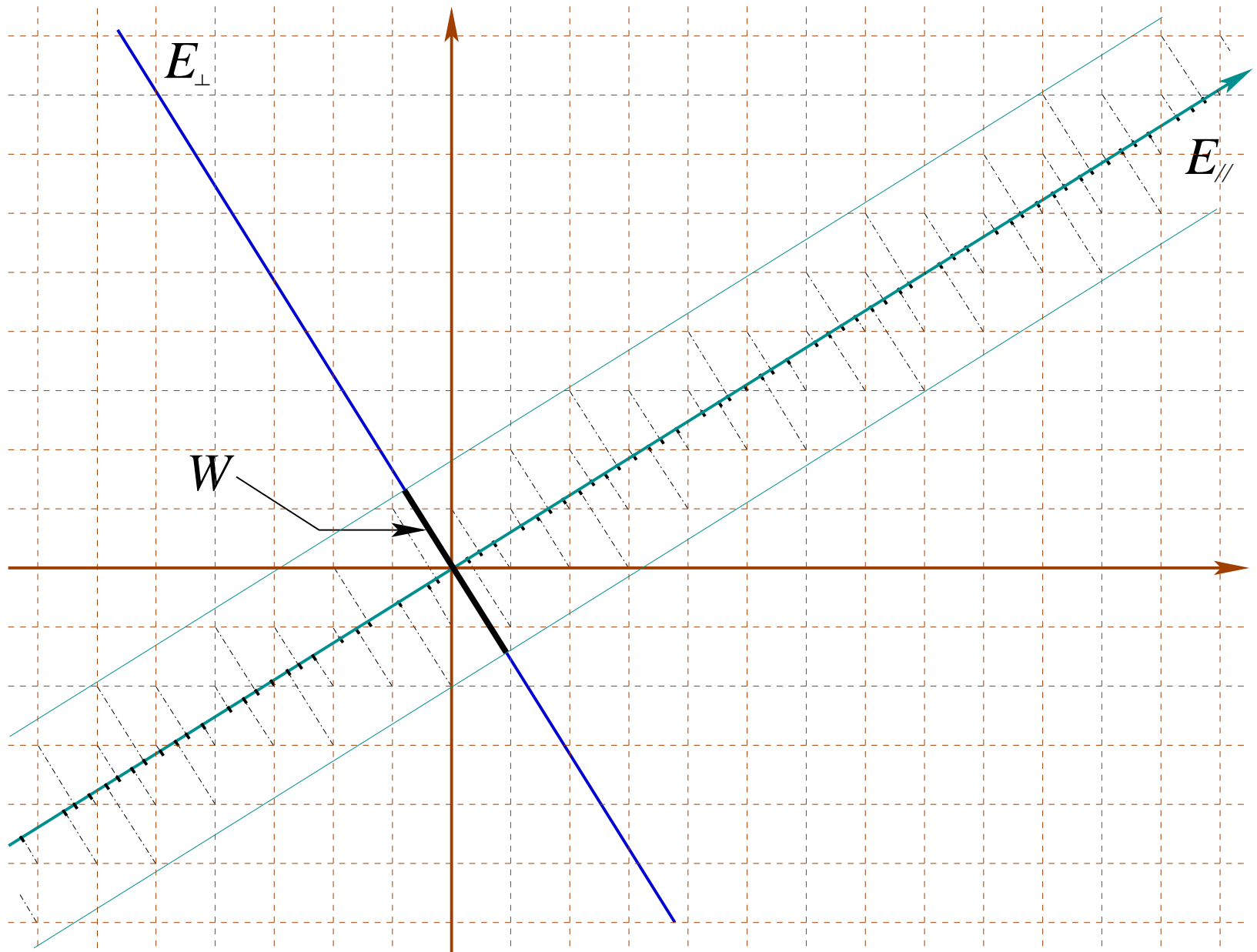
Use the *cut-and-project* construction:

$$\mathbb{R}^d \simeq \mathcal{E}_{\parallel} \xleftarrow{\pi_{\parallel}} \mathbb{R}^n \xrightarrow{\pi_{\perp}} \mathcal{E}_{\perp} \simeq \mathbb{R}^{n-d}$$

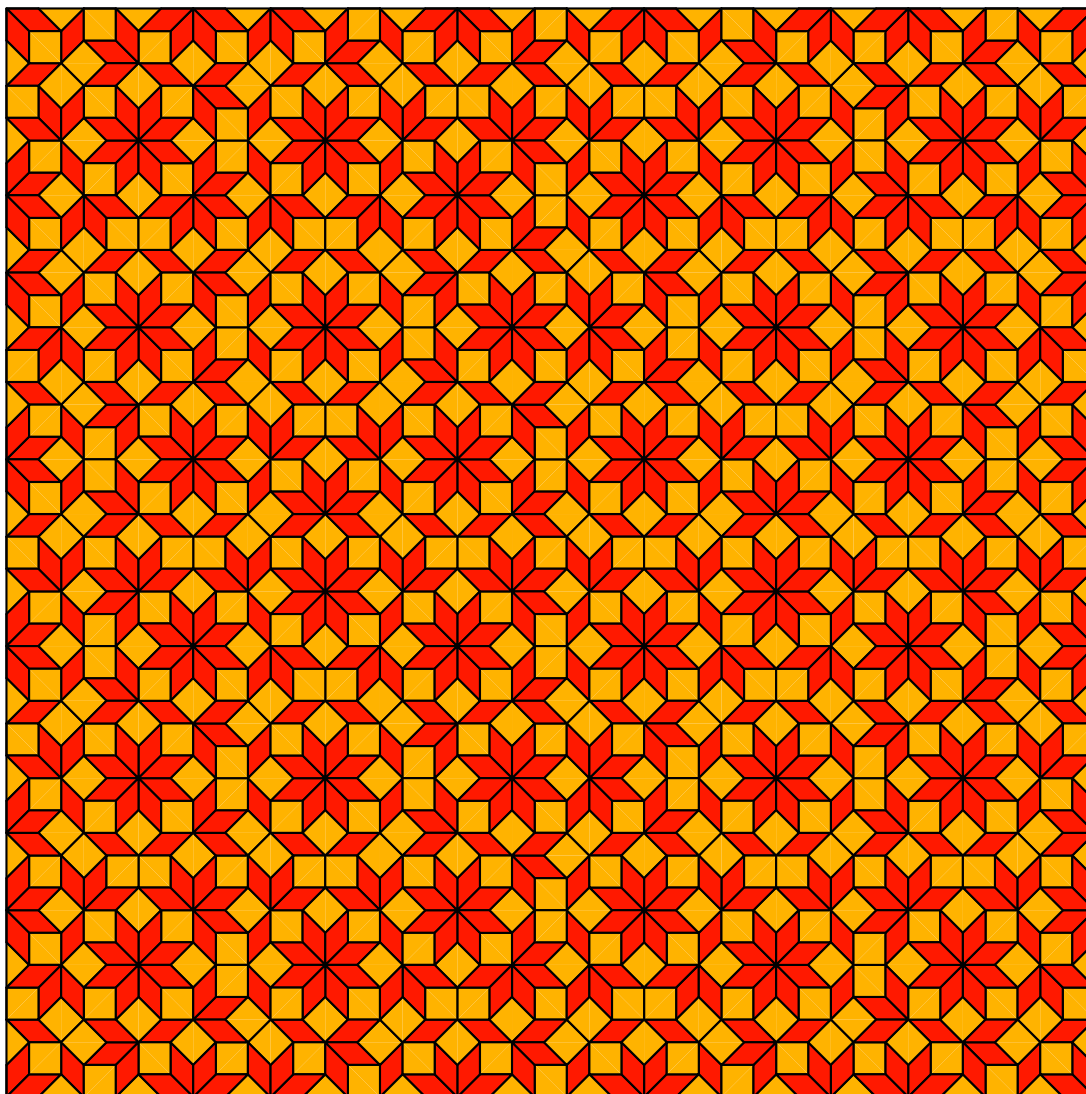
$$\mathcal{L} \xleftarrow{\pi_{\parallel}} \tilde{\mathcal{L}} \xrightarrow{\pi_{\perp}} W ,$$

Here

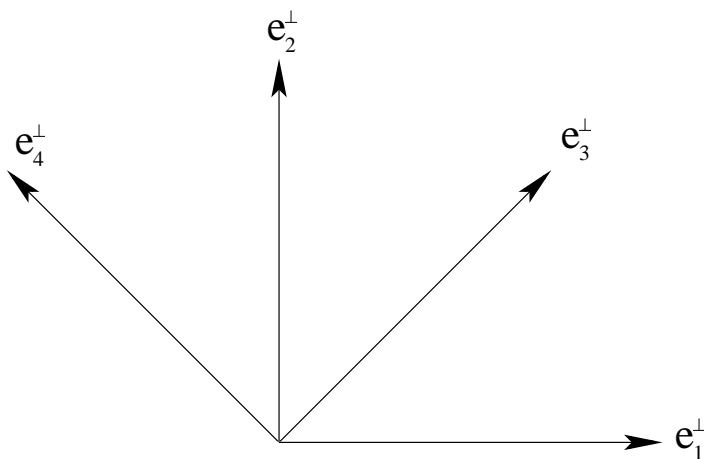
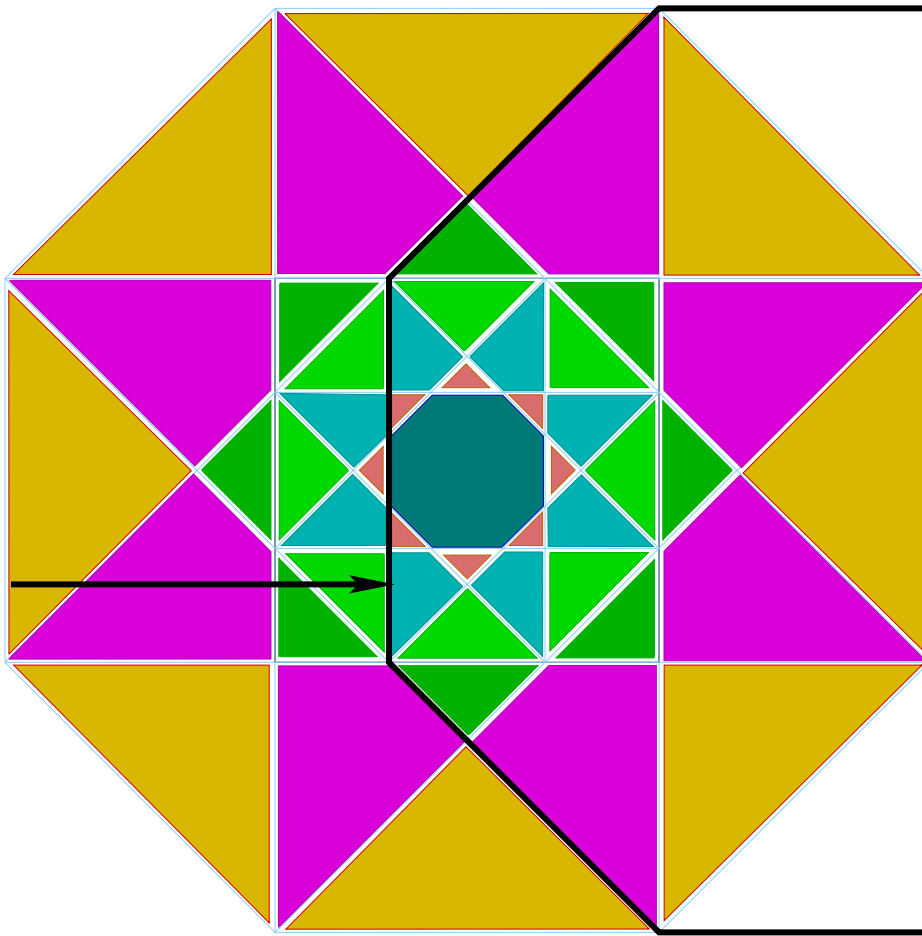
1. $\tilde{\mathcal{L}}$ is a *lattice* in \mathbb{R}^n ,
2. the *window* W is a compact polytope.
3. \mathcal{L} is the induced *quasilattice* in \mathcal{E}_{\parallel} .



– The cut-and-project construction –



- The octagonal tiling -



- The transversal of the Octagonal Tiling -
- is completely disconnected -

II.3)- Hull of Quasicrystals

1. Let \mathcal{F} be the family of affine hyperplanes in \mathbb{R}^n with projections on \mathcal{E}_\perp containing the maximal faces of W . Endow \mathbb{R}^n with the topology such that for any $F \in \mathcal{F}$ and $a \in \tilde{\mathcal{L}}$, the two half spaces separated by the affine hyperplane $F + a$ are both closed and open. Let $\mathbb{R}_{\mathcal{F}}^n$ be the completion of \mathbb{R}^n with this topology.
2. By construction, for $a \in \tilde{\mathcal{L}}$, the map $\tilde{x} \in \mathbb{R}^n \mapsto \tilde{x} + a \in \mathbb{R}^n$ extends to $\mathbb{R}_{\mathcal{F}}^n$ by continuity. Let then $\mathbf{T}_{\mathcal{F}}^n = \mathbb{R}_{\mathcal{F}}^n / \tilde{\mathcal{L}}$ be the corresponding *pseudo torus*.
3. By construction, for $x \in \mathcal{E}_\parallel$, the map $\tilde{x} \in \mathbb{R}^n \mapsto \tilde{x} + x \in \mathbb{R}^n$ extends also by continuity to $\mathbb{R}_{\mathcal{F}}^n$ and commutes with $\tilde{\mathcal{L}}$. Thus it defines an \mathbb{R}^d action on $\mathbf{T}_{\mathcal{F}}^n$.

Theorem 1 \mathcal{L} is a Meyer set. Its Hull is conjugate to $\mathbf{T}_{\mathcal{F}}^n$ endowed with its canonical \mathbb{R}^d -action. This Hull is uniquely ergodic.

III - Tilings & Point Sets

C. RADIN, *Miles of Tiles*, Student Mathematical Library, Vol 1, Amer. Math. Soc., Providence, (1999).

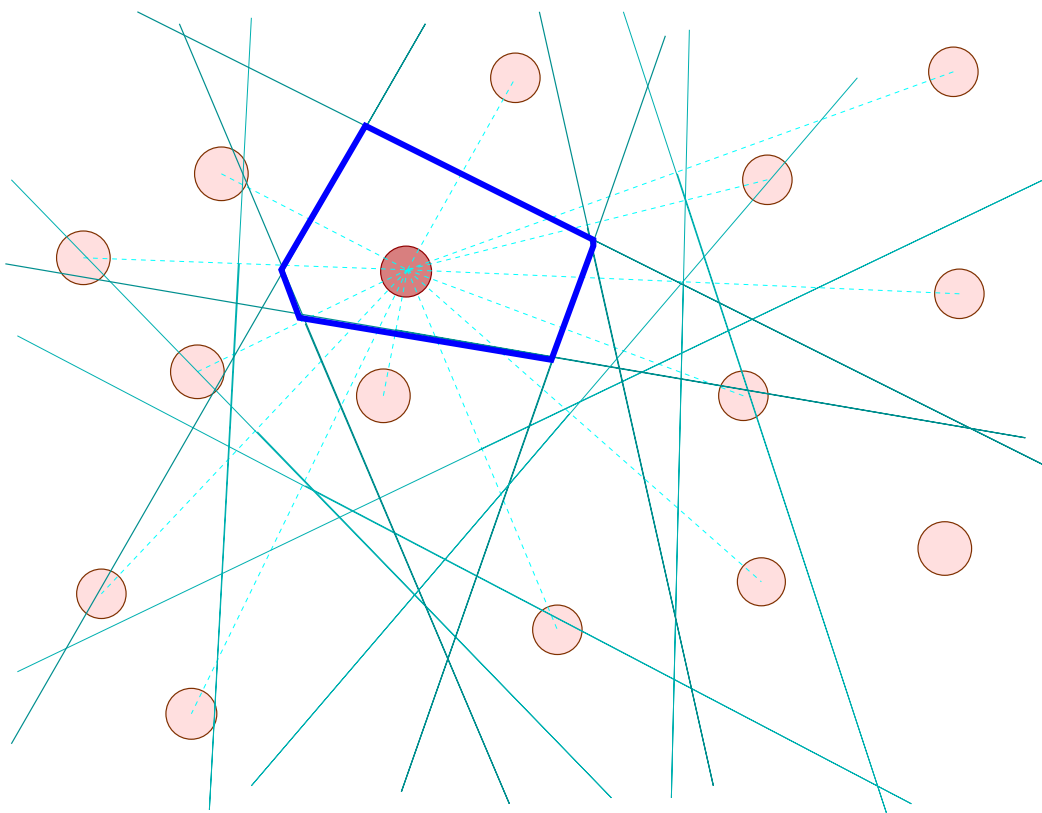
M. SENECHAL, *Crystalline Symmetries : An informal mathematical introduction*, Institute of Physics.
Alan Hilger, Ltd., (1990).

J. BELLISSARD, R. BENEDETTI, J.-M. GAMBAUDO, *Spaces of Tilings, Finite Telescopic Approximations and Gap-Labeling*,
preprint August (2001).

III.1)- Voronoi Cells

For \mathcal{L} Delone and $x \in \mathcal{L}$, the *Voronoi cell of x* is

$$V_x = \{y \in \mathbb{R}^d; |y - x| < |y - x'|, \forall x' \in \mathcal{L} \setminus \{x\}\}$$



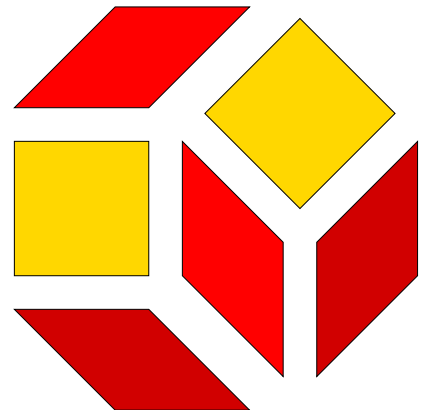
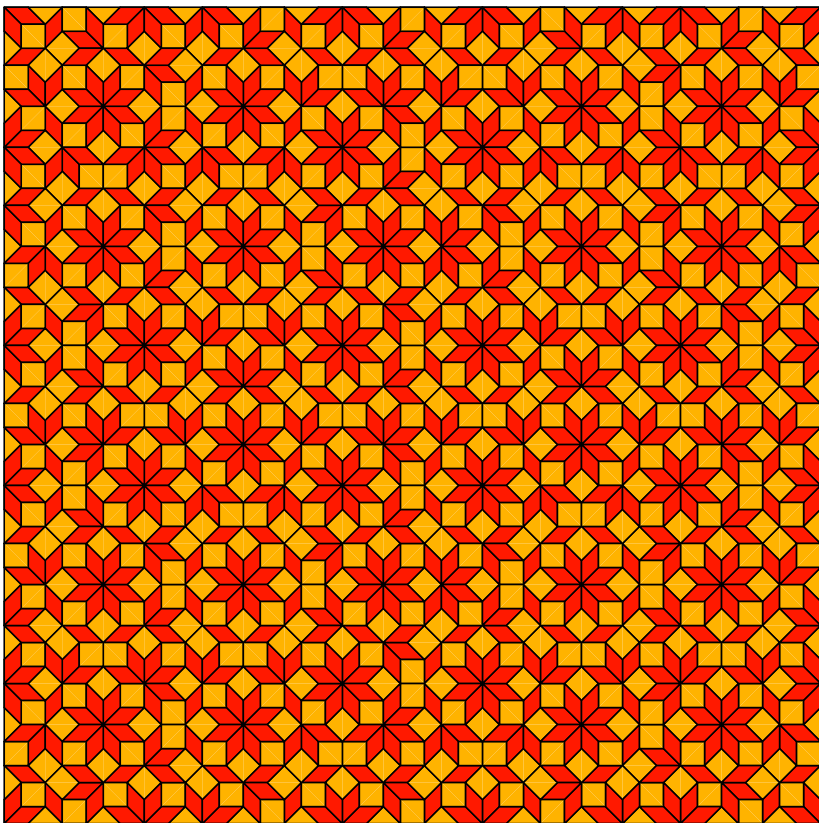
- *Building a Voronoi Cell* -

The V_x 's are open polyhedrons with uniformly bounded diameter. They are mutually disjoint and their closure cover \mathbb{R}^d : it is a *tiling of \mathbb{R}^d*

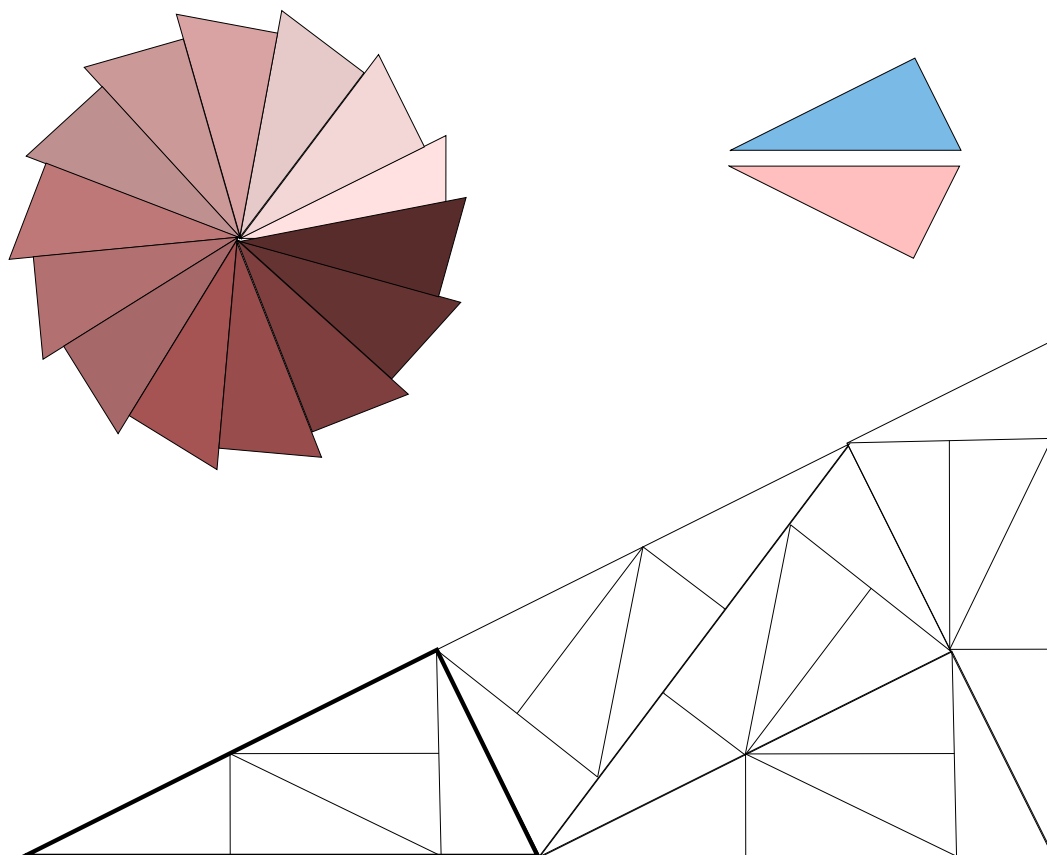
III.2)- The Finite Pattern Condition

A *patch* is the set of tiles of \mathcal{T} contained in some ball. A tiling \mathcal{T} fulfills the *finite pattern condition* (FPC) if the number of patches of radius smaller than R *modulo translations* is finite for all R 's.

Then the transversal is a Cantor set.



- *The octagonal tiling is FPC* -

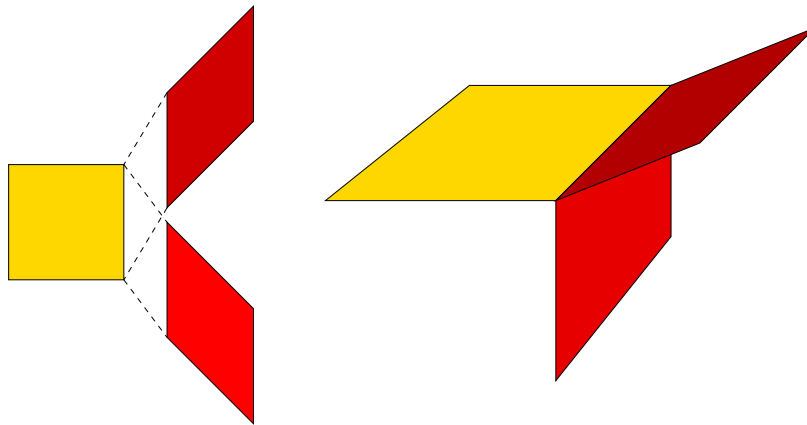


- *The pinwheel tiling is NOT FPC !* -

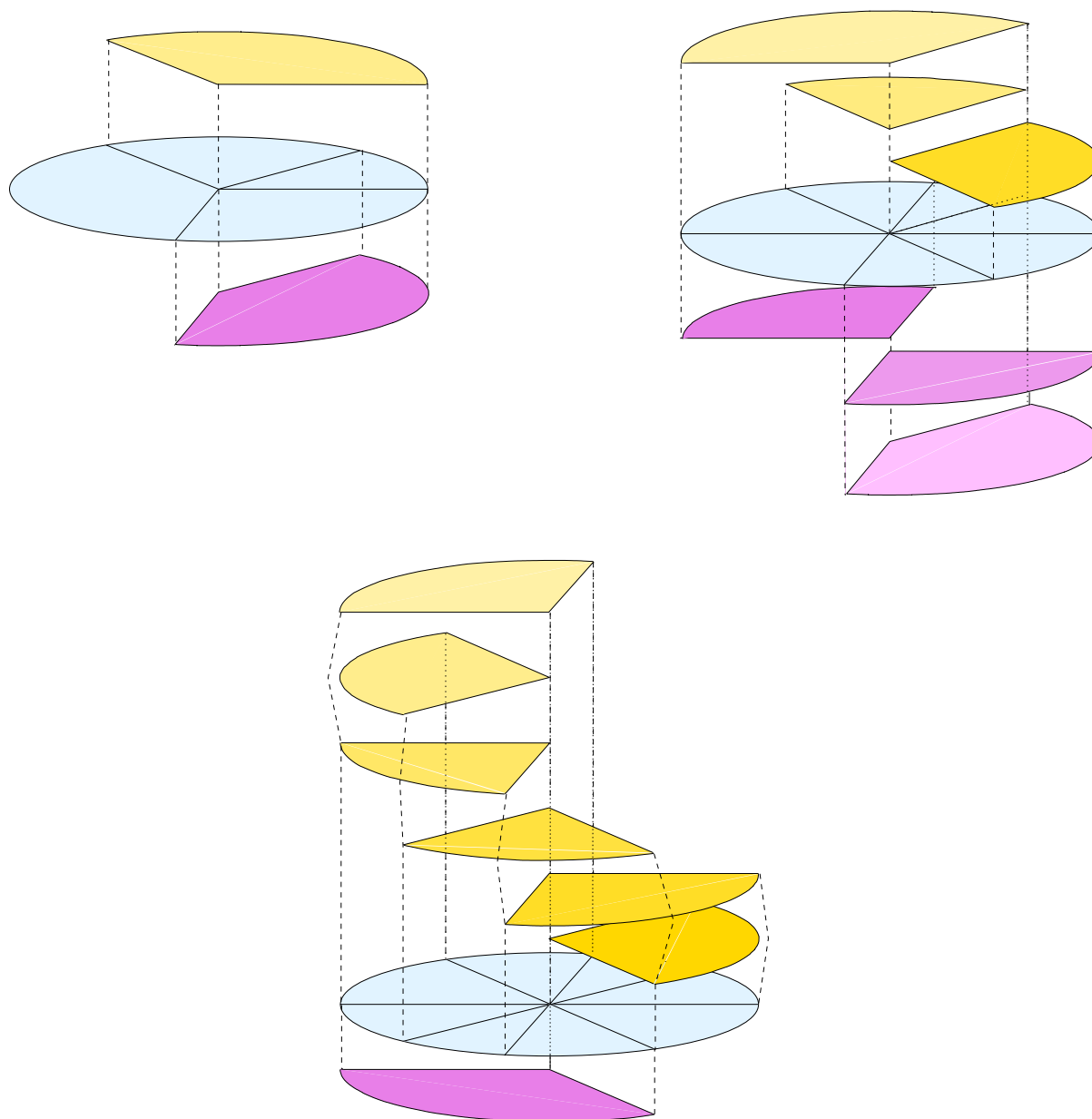
III.3)- Branched Oriented Flat Manifolds

Step 1:

1. X is the disjoint union of all *prototiles*;
2. glue prototiles T_1 and T_2 along a face $F_1 \subset T_1$ and $F_2 \subset T_2$ if F_2 is a translated of F_1 and if there are $x_1, x_2 \in \mathbb{R}^d$ such that $x_i + T_i$ are tiles of \mathcal{T} with $(x_1 + T_1) \cap (x_2 + T_2) = x_1 + F_1 = x_2 + F_2$;
3. after identification of faces, X becomes a *branched oriented flat manifold* (BOF) B_0 .



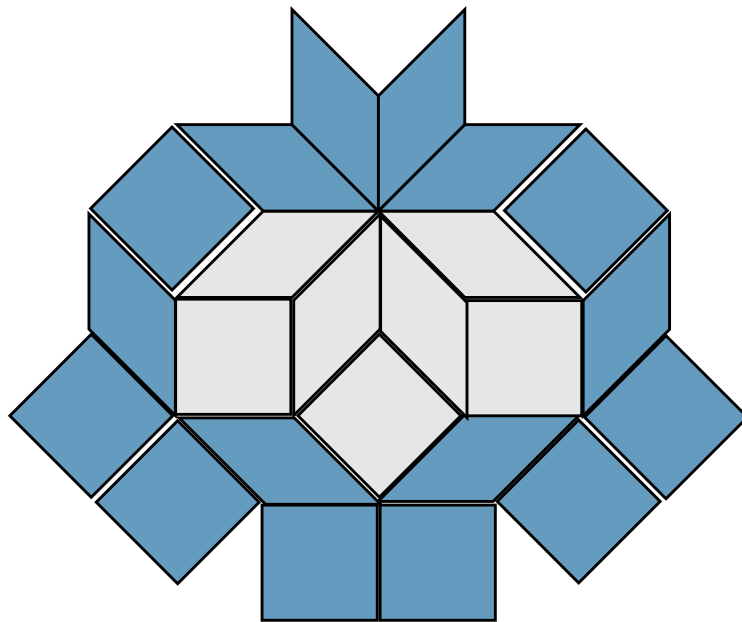
- *The branching process* -



- *Vertex branching for the octagonal tiling* -

Step 2:

1. Choose an increasing sequence $\{R_n\}_{n>0}$ of positive real numbers with $R_n \uparrow \infty$;
2. for each $n \geq 1$ consider all patches of diameter less than R_n ;
3. add to each patch in \mathcal{T} , the tiles touching it from outside along its frontier. Call such a patch *modulo translation* a *colored patch*;
4. proceed then as in Step 1 by replacing prototiles by colored patches to get the BOF-manifold B_n .



- A colored patch -

Step 3:

1. Define a *BOF-submersion* $f_n : B_{n+1} \mapsto B_n$ by identifying patches of order n in B_{n+1} with the protiles of B_n ;
2. call Ω the *projective limit* of the sequence

$$\dots \xrightarrow{f_{n+1}} B_{n+1} \xrightarrow{f_n} B_n \xrightarrow{f_{n-1}} \dots$$

3. there are commuting vector fields X_1, \dots, X_d on B_n generating local translations and giving rise to a \mathbb{R}^d action \mathbb{T} on Ω .

Theorem 2 *The dynamical system*

$$(\Omega, \mathbb{R}^d, \mathbb{T}) = \varprojlim (B_n, f_n)$$

obtained as inverse limit of branched oriented flat manifolds, is conjugate to the Hull of the Delone set of the tiling \mathcal{T} by an homeomorphism.

IV - NC Brillouin Zone

J. BELLISSARD, *The Gap Labelling Theorems for Schrödinger's Operators*,
in *From Number Theory to Physics*, pp. 538-630, Les Houches March 89, Springer,
J.M. Luck, P. Moussa & M. Waldschmidt Eds., (1993).

IV.1)- Algebra

$(\Omega, \mathbb{R}^d, \tau)$ is a topological dynamical system. One orbit at least is dense. The crossed product

$$\mathcal{A} = \mathcal{C}(\Omega) \rtimes_{\tau} \mathbb{R}^d$$

is (almost) the smallest C^* -algebra containing both the space of continuous functions on Ω and the action of \mathbb{R}^d submitted to the commutation rules (for $f \in \mathcal{C}(\Omega)$)

$$T(a) f T(a)^{-1} = f \circ \tau^{-a}$$

1. For a crystal $\Omega = \mathbb{V}$, \mathbb{R}^d acts by quotient action.
2. $\mathcal{C}(\mathbb{V}) \rtimes_{\tau} \mathbb{R}^d \simeq \mathcal{C}(\mathbb{B}) \otimes \mathcal{K}$, where \mathcal{K} is the algebra of compact operators and \mathbb{B} is the dual of the period group of \mathcal{L} . \mathbb{B} is called the *Brillouin zone*.

\mathcal{A} is the *Noncommutative version of the space of \mathcal{K} -valued function over the Brillouin zone.*

IV.2)- Construction of \mathcal{A}

Endow $\mathcal{A}_0 = \mathcal{C}_c(\Omega \times \mathbb{R}^d)$ with (here $A, B \in \mathcal{A}_0$):

1. Product

$$A \cdot B(\omega, x) = \int_{y \in \mathbb{R}^d} d^d y A(\omega, y) B(\tau^{-y} \omega, x - y)$$

2. Involution

$$A^*(\omega, x) = \overline{A(\tau^{-x} \omega, -x)}$$

3. A faithful family of representations in $\mathcal{H} = L^2(\mathbb{R}^d)$

$$\pi_\omega(A) \psi(x) = \int_{\mathbb{R}^d} d^d y A(\tau^{-x} \omega, y - x) \cdot \psi(y)$$

if $A \in \mathcal{A}_0$, $\psi \in \mathcal{H}$.

4. C^* -norm

$$\|A\| = \sup_{\omega \in \Omega} \|\pi_\omega(A)\| .$$

Definition 1 *The C^* -algebra \mathcal{A} is the completion of \mathcal{A}_0 under this norm.*

IV.3)- Calculus

Integration: Let \mathbb{P} be an \mathbb{R}^d -invariant ergodic probability measure on Ω . Then set (for $A \in \mathcal{A}_0$):

$$\mathcal{T}_{\mathbb{P}}(A) = \int_{\Omega} d\mathbb{P} A(\omega, 0) = \overline{\langle 0 | \pi_{\omega}(A) 0 \rangle}^{dis.}$$

Then $\mathcal{T}_{\mathbb{P}}$ extends as a *positive trace* on \mathcal{A} .

Trace per unit volume: thanks to Birkhoff's theorem:

$$\mathcal{T}_{\mathbb{P}}(A) = \lim_{\Lambda \uparrow \mathbb{R}^d} \frac{1}{|\Lambda|} \text{Tr}(\pi_{\omega}(A) \upharpoonright_{\Lambda}) \quad \text{a.e. } \omega$$

Differential calculus: A commuting set of $*$ -derivations is given by

$$\partial_i A(\omega, x) = \imath x_i A(\omega, x)$$

defined on \mathcal{A}_0 . Then $\pi_{\omega}(\partial_i A) = -\imath [X_i, \pi_{\omega}(A)]$ where $X = (X_1, \dots, X_d)$ are the coordinates of the position operator.

IV.4)- Electrons

Schrödinger's equation (ignoring interactions) on \mathbb{R}^d

$$H_\omega = -\frac{\hbar^2}{2m}\Delta + \sum_{y \in \mathcal{L}_\omega} v(\cdot - y) ,$$

acting on $\mathcal{H} = L^2(\mathbf{R}^d)$. Here $v \in L^1(\mathbb{R}^d)$ is real valued, decays fast enough, is the *atomic potential*.

Lattice case (*tight binding representation*)

$$\tilde{H}_\omega \psi(x) = \sum_{y \in \mathcal{L}_\omega} h(\mathbb{T}^{-x}\omega, y - x) \psi(y) ,$$

Proposition 3 1. There is $R(z) \in \mathcal{A}$, such that, for every $\omega \in \Omega$ and $z \in \mathbb{C} \setminus \mathbb{R}$

$$(z - H_\omega)^{-1} = \pi_\omega(R(z)) .$$

2. There is $\tilde{H} \in C^*(\Gamma_{tr})$ such that $\tilde{H}_\omega = \pi_\omega(\tilde{H})$.

3. If $\Sigma_H = \bigcup_{\omega \in \Omega} \text{Sp}(H_\omega)$, then $R(z)$ is holomorphic in $z \in \mathbb{C} \setminus \Sigma_H$. The bounded components of $\mathbb{R} \setminus \Sigma_H$ are called *spectral gaps* (same with \tilde{H}).

IV.5)- Density of States

- Let \mathbb{P} be an invariant ergodic probability on Ω . Let

$$\mathcal{N}(E) = \lim_{\Lambda \uparrow \mathbf{R}^d} \frac{1}{|\Lambda|} \# \{ \text{eigenvalues of } H_\omega|_\Lambda \leq E \}$$

It is called the *Integrated Density of states* or *IDS*.

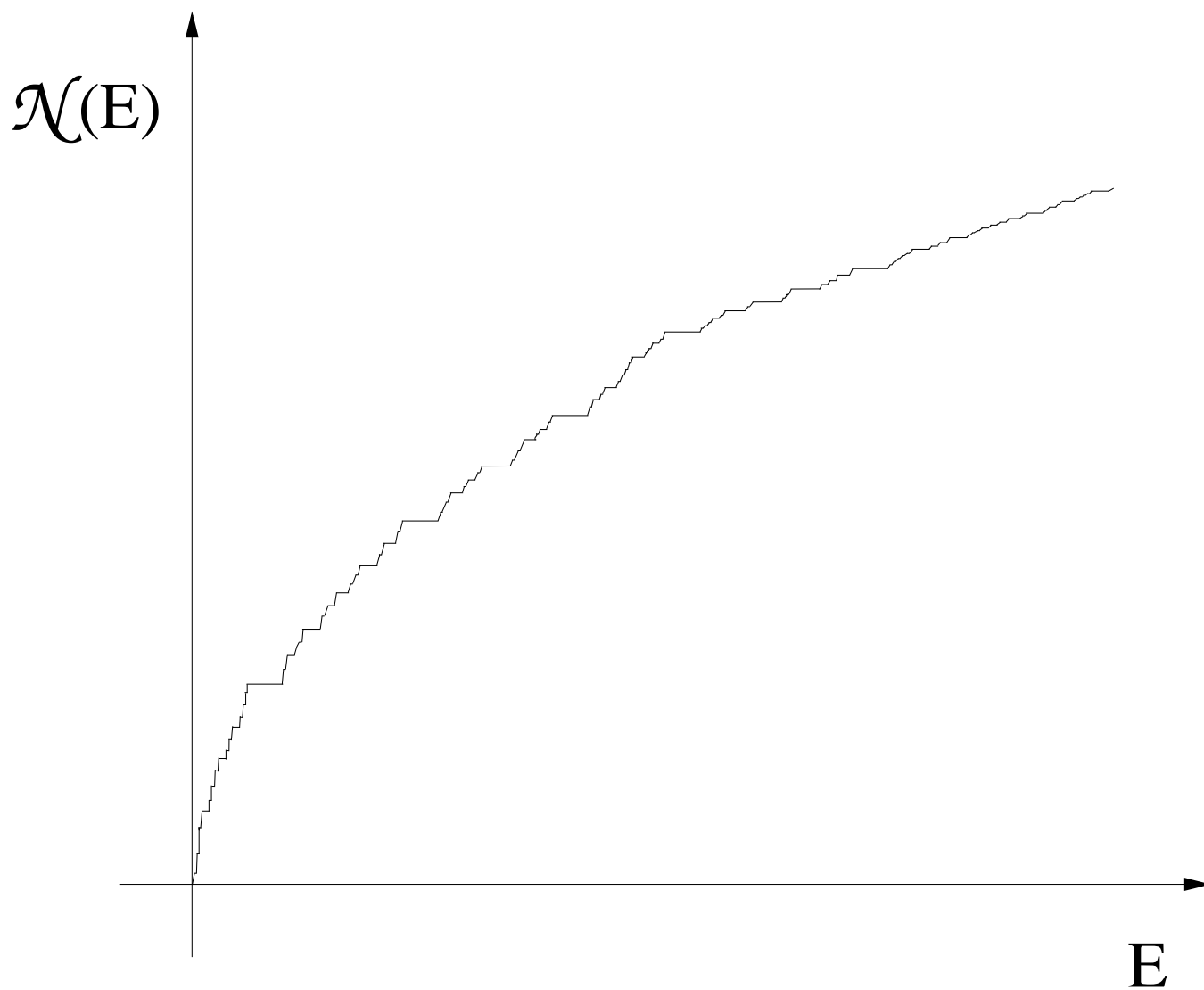
- The limit above exists \mathbb{P} -almost surely and

$$\mathcal{N}(E) = \mathcal{T}_{\mathbb{P}}(\chi(H \leq E)) \quad (\text{Shubin, '76})$$

$\chi(H \leq E)$ is the eigenprojector of H in $\mathcal{L}^\infty(\mathcal{A})$.

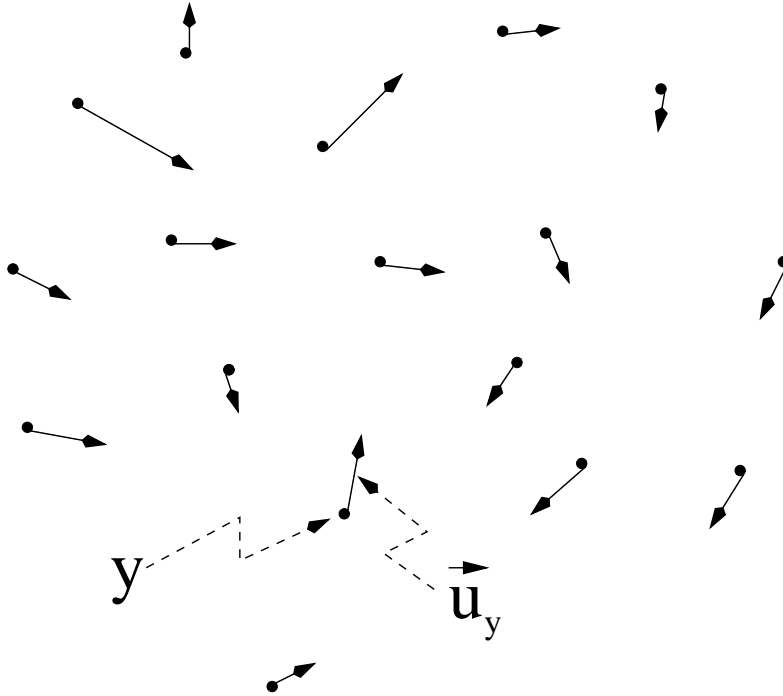
- \mathcal{N} is non decreasing, non negative and constant on gaps. $\mathcal{N}(E) = 0$ for $E < \inf \Sigma_H$. For $E \rightarrow \infty$, $\mathcal{N}(E) \sim \mathcal{N}_0(E)$ where \mathcal{N}_0 is the IDS of the free case (namely $v = 0$).

- *Gaps can be labelled by the value the IDS takes on them*



- An example of IDS -

IV.6)- Phonons



1. Phonons are *acoustic waves* produced by small displacements of the atomic nuclei.
2. These waves are polarized with d -directions of polarization: $d - 1$ are *transverse*, one is *longitudinal*.
3. The nuclei motion is approximatively harmonic and quantized according to the *Bose-Einstein* statistics.
4. The charged nuclei interact with electrons, leading to an *electron-phonon interaction*.

1. For identical atoms with *harmonic motion*, the classical equations of motion are:

$$M \frac{d^2 \vec{u}_{(\omega, x)}}{dt^2} = \sum_{x \neq y \in \mathcal{L}_\omega} K_\omega(x, y) (\vec{u}_{(\omega, y)} - \vec{u}_{(\omega, x)})$$

where M is the atomic mass, $\vec{u}_{(\omega, x)}$ is its classical displacement vector and $K_\omega(x, y)$ is the matrix of *spring constants*.

2. $K_\omega(x, y)$ decays fast in $x - y$, uniformly in ω ..
3. Covariance gives

$$K_\omega(x, y) = k(\tau^{-x}\omega, y - x)$$

thus

$$k \in C^*(\Gamma_{tr}) \otimes M_d(\mathbb{C})$$

4. Then the spectrum of k/M gives the *eigenmodes* propagating in the solid. Its density (DPM) is given by Shubin's formula again.

V - Gap labels and K -theory

B. BLACKADAR, *K-Theory for Operator Algebras*,
Springer, New-York, (1986), 2nd Ed, Cambridge U. Press (1998)

J. BELLISSARD, *The Gap Labelling Theorems for Schrödinger's Operators*,
in *From Number Theory to Physics*, pp. 538-630, Les Houches March 89, Springer,
J.M. Luck, P. Moussa & M. Waldschmidt Eds., (1993).

J. BELLISSARD, R. BENEDETTI, J. M. GAMBAUDO, *Spaces of Tilings,
Finite Telescopic Approximations and Gap-Labeling*,
math.DS/0109062, (2001).

M. BENAMEUR, H. OYONO-OYONO, *Gap labeling for quasicrystals
(proving a conjecture of J. Bellissard)*,
math.KT/0112113, (2001).

J. KAMINKER, I.F. PUTNAM, *A proof of the gap labeling conjecture*,
math.KT/0205102, (2002).

V.1)- K -group

Let \mathcal{A} be a separable C^* -algebra .

1. A *projection* is $P \in \mathcal{A}$ with $P = P^* = P^2$.
2. Two projections P and Q are *equivalent* if there is $U \in \mathcal{A}$ such that $P = UU^*$, $Q = U^*U$.
 $[P]$ denotes the equivalent classe of P .
3. Two projections P and Q are *orthogonal* if $PQ = QP = 0$. Then $P + Q$ is a projection called the *direct sum* $P \oplus Q$.
4. If $\mathcal{K} = \overline{\bigcup_n M_n(\mathbb{C})}^{\|\cdot\|}$, then \mathcal{A} is *stable* if $\mathcal{A} \simeq \mathcal{A} \otimes \mathcal{K}$. In a stable C^* -algebra for any pair P, Q of projections, there are $P' \sim P$, $Q' \sim Q$ with $P'Q' = Q'P' = 0$.
 Thus $[P] + [Q] = [P' \oplus Q']$ is well defined.
5. $K_0(\mathcal{A})$ is the group generated by the $[P]$'s with the previous addition.
6. $[P]$ is invariant by homotopy.

V.2)- K -group labels

- If E belongs to a gap \mathfrak{g} , the characteristic function $E' \in \mathbf{R} \mapsto \chi(E' \leq E)$ is continuous on the spectrum of H . Thus:

$P_{\mathfrak{g}} = \chi(H \leq E)$ is a projection in \mathcal{A} !

- $\mathcal{N}(E) = \mathcal{T}_{\mathbb{P}}(P_{\mathfrak{g}}) \in \mathcal{T}_{\mathbb{P}}^*(K_0(\mathcal{A}))$!

Theorem 3 (Abstract gap labeling theorem)

- $S \subset \Sigma_H$ clopen, $n_S = [\chi_S(H)] \in K_0(\mathcal{A})$. If $S_1 \cap S_2 = \emptyset$ then $n_{S_1 \cup S_2} = n_{S_1} + n_{S_2}$ (**additivity**).
- Gap labels are invariant under norm continuous variation of H (**homotopy invariance**).
- For $\lambda \in [0, 1] \mapsto H(\lambda) \in \mathcal{A}$ continuous, if $S(\lambda) \subset \Sigma_H$ clopen, continuous in λ with $S(0) = S_1 \cup S_2$, $S(1) = S'_1 \cup S'_2$ and $S_1 \cap S_2 = \emptyset = S'_1 \cap S'_2$ then $n_{S_1} + n_{S_2} = n_{S'_1} + n_{S'_2}$ (**conservation of gap labels under band crossings**).

Theorem 4 *If \mathcal{L} is an finite type Delone set in \mathbb{R}^d with Hull $(\Omega, \mathbb{R}^d, \mathbb{T})$, then, for any \mathbb{R}^d -invariant probability measure \mathbb{P} on Ω*

$$\mathcal{T}_{\mathbb{P}}^* (K_0(\mathcal{A})) = \int_X d\mathbb{P}_{tr} \mathcal{C}(X, \mathbb{Z}) .$$

if $\mathcal{A} = \mathcal{C}(\Omega) \rtimes \mathbb{R}^d$, X is the canonical transversal and \mathbb{P}_{tr} the transverse measure induced by \mathbb{P} .

Main ingredient for the proof

*the Connes measured index theorem
for foliated spaces*

A. CONNES, *Sur la théorie non commutative de l'intégration*,
Lecture Notes in Math **725**, 19-143, Springer, Berlin (1979).

V.3)- History

For $d = 1$ this result follows from the Pimsner & Voiculescu exact sequence (*Bellissard, '92*).

For $d = 2$, a double use of the Pimsner & Voiculescu exact sequence provides the result (*van Elst, '95*).

For $d \geq 3$ whenever $(\Omega, \mathbb{R}^d, \mathbb{T})$ is Morita equivalent to a \mathbb{Z}^d -action, using spectral sequences (*Hunton, Forrest*) this theorem was proved for $d = 3$ (*Bellissard, Kellendonk, Legrand, '00*).

The theorem has also been proved for all d 's recently and independently by

- (*Bellissard, Benedetti, Gambaudo, 2001*),
- (*Benamou, Oyono, 2001*),
- (*Kaminker, Putnam, 2002*).